

1 **Long-term global ocean heat content change driven by sub-polar**
2 **surface heat fluxes**

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7 **Key Points:**

- 8 • We introduce a novel tracer-percentile framework which relates ocean heat content
9 trends to surface flux and turbulent mixing changes
10 • Heat content changes in the 90% coldest ocean volume are traced to heat fluxes
11 which, on average, enter 23% of the surface area of the ocean
12 • Using this framework, we trace a cooling bias in the 5-20% warmest volume of
13 CMIP6 climate models to surface flux biases

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Abstract

The ocean has absorbed approximately 90% of the accumulated heat in the climate system since 1970. As global warming accelerates, understanding ocean heat content changes and tracing these to surface heat input is becoming increasingly important. We introduce a novel tracer-percentile framework in which we organise the ocean into temperature percentiles from warmest to coldest, allowing us to trace changes in ocean temperature to changes in air-sea heat fluxes and mixing. Applying this framework to observations and historical CMIP6 simulations, we find that 40% of heat uptake between 1970 and 2014 occurs in the warmest 10% ocean volume. However, the coolest 90% ocean volume outcrops over 23% of the ocean’s surface area, implying a disproportionately large heat input per unit area. Additionally, a cold bias in the CMIP6 models is traced to inaccurate sea surface temperatures and surface heat fluxes into the warmest 5 – 20% ocean volume.

Plain Language Summary

The ocean has absorbed approximately 90% of the heat that has built up in the climate system since 1970. Understanding the processes which drive this uptake of heat by the ocean is critical to climate projections. Typically, this has required the use of climate and ocean models, which must be validated with ocean observations. To enable this, we introduce a new method, the tracer-percentile framework, which allows us to directly use observations to understand the processes dictating heat uptake by the ocean. Using climate models (collectively called CMIP6) and ocean observations, we calculate the heat input into layers of water at the ocean surface, the total heat stored in these layers, and the mixing of heat between these layers. We find that the heat input at the surface and heat stored in the ocean has increased between 1970 and 2014, with 60% of the increase in heat uptake at the surface happening over about a quarter of the ocean’s surface, which connects to 90% of the world’s ocean volume. We also identify inaccuracies in the CMIP6 models and trace these problems to the way the surface properties of the ocean (or surface heat input) are represented.

1 Introduction

The global ocean has absorbed approximately 90% of the excess heat in the climate system since 1970 [Schuckmann *et al.*, 2020], impacting global sea level rise [Domingues *et al.*, 2008; Church and White, 2011; Gregory *et al.*, 2013], surface air temperature [Watanabe *et al.*, 2013] and extreme weather [Lin *et al.*, 2013]. As the radiative imbalance of the planet is so difficult to measure directly, knowledge of global ocean heat content (OHC) changes resulting from ocean heat uptake is critical for monitoring climate change. Heat enters the ocean at the surface and is subsequently advected and diffused in the ocean [Gregory, 2000; Kuhlbrodt *et al.*, 2015; Liang *et al.*, 2015; Cummins *et al.*, 2016]. Past research has thus articulated a *fixed-depth framework*, which quantifies changes in global OHC at a given depth as a balance between surface fluxes and vertical advection and diffusion. Analysing ocean circulation and heat transport in fixed-depth coordinates has yielded a number of insights. Prior work has established that mean downward heat transport is balanced by along-isopycnal upward eddy fluxes, particularly in the Southern Ocean [Gregory, 2000; Gnanadesikan *et al.*, 2005; Wolfe *et al.*, 2008; Morrison *et al.*, 2013; Kuhlbrodt *et al.*, 2015]. In addition, Zika *et al.* [2013] found that downward heat transport in the meridional overturning circulation is driven by a combination of salinification in the subtropics and wind-driven Ekman pumping in the Southern Ocean.

While the fixed-depth framework provides a useful way of analysing heat transport, it can be challenging to *quantify* observed OHC, as this requires knowledge of the velocity field [something that is only possible with models (e.g., Wolfe *et al.* [2008]; Kuhlbrodt

63 *et al.* [2015]) and reanalysis products (e.g., *Liang et al.* [2015, 2017])). Importantly, the
 64 vertical heat budget is impacted by adiabatic processes, internal variability and heave,
 65 meaning that adiabatic redistribution of existing heat impacts trends in OHC at fixed depth.
 66 Therefore, there remains scope to develop a framework which quantifies how much ob-
 67 served heat is added to the climate system at a given location.

68 Numerous studies have moved towards analysing ocean circulation and heat transport
 69 in a water mass-based reference frame [*Walín, 1982; Groeskamp et al., 2019*]. Recently
 70 this framework has been used to understand climate change [*Zika et al., 2015*] and climate
 71 variability [*Evans et al., 2017*]. *Holmes et al.* [2019] formalised a heat budget in fixed-
 72 temperature coordinates, termed the diathermal heat transport framework. By analysing
 73 heat transport across a given *isotherm*, adiabatic processes, internal variability and heave
 74 are excluded. In the resulting *temperature-based framework*, heat content tendencies are
 75 tracked to changes in diabatic transport processes - air-sea fluxes and mixing [*Holmes*
 76 *et al., 2019*]. The temperature-based framework allows us to link the temperature at which
 77 heat enters the ocean to the temperature classes which exhibit OHC changes. Diathermal
 78 fluxes may also be calculated purely from observed in-situ temperature and surface flux
 79 data (with the mixing term calculated by residual).

80 The temperature-based framework addresses the challenges posed by the fixed-depth
 81 framework, but also presents its own complications. For instance, isotherms shift as the
 82 ocean warms such that the characteristic isotherms of a given region move to a different
 83 region, making partitioning between the tropics, sub-tropics and sub-polar oceans dif-
 84 ficult. We introduce a novel diagnostic which defines OHC changes at temperature per-
 85 centiles, ordered from warmest to coldest. The new *tracer-percentile framework* allows us
 86 to quantify observed changes in OHC and trace them to changes in surface heat fluxes and
 87 mixing using a coarse-resolution hydrographic observational dataset. The tracer-percentile
 88 framework also enables a direct evaluation of model biases and attribution of biases to
 89 surface fluxes or mixing. In a uniformly warming ocean, the tracer-percentile framework
 90 avoids issues associated with shifting isotherms. In addition, organising the ocean by vol-
 91 ume means that OHC changes can be translated into changes in the temperature of the
 92 surface that bounds a given volume percentile, allowing a comparison with changes in the
 93 globally-averaged temperature at fixed depth.

94 In section 2, we define the heat budget in terms of the tracer-percentile framework.
 95 In section 3, we summarise the data sources used in the analysis. In section 4, we detail
 96 the key results of the study, and in section 5 we summarise this work and discuss impor-
 97 tant implications.

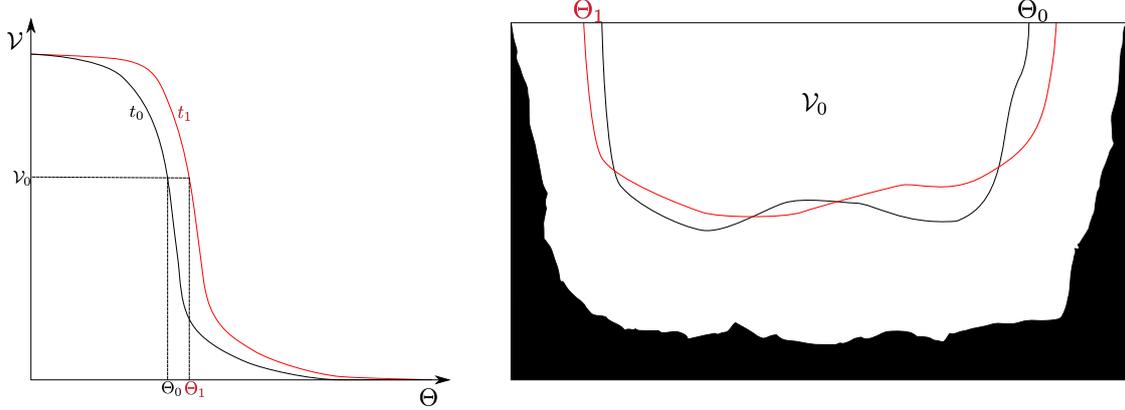
98 2 Theory

99 2.1 The tracer-percentile framework

100 In the fixed-depth framework, the variable of interest is $H(z, t)$, the total OHC above
 101 a given depth z (equivalent to the volume-integral of the temperature multiplied by the
 102 reference density ρ_0 and specific heat capacity of seawater C_p , in the Boussinesq case).
 103 The fixed-temperature framework, on the other hand, considers the total OHC above a
 104 given isotherm $\mathcal{H}(\Theta^*, t)$ where the corresponding volume is denoted $\mathcal{V}(\Theta^*, t)$ [*Palmer*
 105 *and Haines, 2009; Holmes et al., 2019*]:

$$\begin{aligned} \mathcal{V}(\Theta^*, t) &= \iiint_{\Theta(x, y, z, t) > \Theta^*} dx dy dz; \\ \mathcal{H}(\Theta^*, t) &= \rho_0 C_p \int_0^{\mathcal{V}(\Theta^*, t)} \Theta(x, y, z, t) d\mathcal{V}, \end{aligned} \quad (1)$$

106 where Θ is the three-dimensional temperature field and Θ^* is the binning temperature over
 107 which the integration occurs.



108 **Figure 1.** Illustration of the change in bounding temperature of a fixed volume in a warming ocean. *Left:*
 109 Cumulative distribution of volume as a function of temperature, ordered from hot to cold, at times t_0 (black)
 110 and t_1 (red), where $t_1 > t_0$. *Right:* Zonally-averaged representation of the volume \mathcal{V}_0 and the bounding temper-
 111 ature $\Theta_0(\mathcal{V}_0, t_0)$ (black) and $\Theta_1(\mathcal{V}_0, t_1)$ (red).

112 The cumulative distribution function (CDF) of volume $\mathcal{V}(\Theta, t)$ in temperature co-
 113 ordinates (organised from warmest to coldest) is illustrated in figure 1 at times t_0 and t_1
 114 (where the ocean is assumed to warm uniformly for illustrative purposes). As the ocean
 115 warms, assuming the thermal expansion of sea water negligibly modifies the total volume
 116 of the ocean, the CDF moves along the temperature axis while remaining fixed in the cum-
 117 ulative volume axis. Hence, at a fixed temperature the volume increases, and at a fixed
 118 volume the bounding temperature increases. In the *tracer-percentile framework*, we invert
 119 the volume equation in (1) to get the bounding tracer (in this case temperature) $\Theta^{\mathcal{V}}(\mathcal{V}, t)$
 120 of a given volume $\mathcal{V}(\Theta, t)$:

$$\mathcal{V}(\Theta^*, t) \iff \Theta^{\mathcal{V}}(\mathcal{V}, t), \quad (2)$$

121 The temperature-percentile is $p(\Theta^*, t) = 100 \times \mathcal{V}(\Theta^*, t) / \mathcal{V}_T$, where \mathcal{V}_T is the total volume
 122 of the ocean. Note that we ignore the negligible changes in the total volume of the ocean
 123 with time. OHC in temperature percentiles is therefore expressed as:

$$\mathcal{H}(p, t) = 0.01 \mathcal{V}_T \rho_0 C_p \int_0^p \Theta^p(p, t) dp, \quad (3)$$

124 where $\Theta^p(p, t)$ is equivalent to $\Theta^{\mathcal{V}}(\mathcal{V}, t)$ at a given temperature-percentile.

125 2.2 Ocean heat content tendency

126 The OHC tendency in temperature percentiles, $\partial \mathcal{H}(p, t) / \partial t$, may be related to spe-
 127 cific diabatic heat transport processes, providing insight into the drivers of OHC changes.
 128 To formulate the OHC tendency budget in temperature percentiles, we begin by looking
 129 at the diathermal heat transport budget introduced by *Holmes et al.* [2019] for the fixed-
 130 temperature framework.

131 In the fixed-temperature framework, the total heat content tendency $\partial \mathcal{H}(\Theta, t) / \partial t$ is
 132 a consequence of changes in the surface forcing \mathcal{F} , mixing \mathcal{M} and diathermal heat ad-

133 vection $\mathcal{G}\Theta\rho_0C_p$. The diabatic across-isotherm volume flux $\mathcal{G}\Theta\rho_0C_p$ is itself a conse-
 134 quence of surface forcing and mixing. It is also dependent upon an arbitrary choice of
 135 reference temperature Θ_{ref} . *Palmer and Haines* [2009] and *Holmes et al.* [2019] sought to
 136 exclude the reference temperature-dependent diathermal advection term from the heat bud-
 137 get, *Palmer and Haines* [2009] by setting reference temperature as the mean temperature
 138 of the volume bounded by an isotherm, and *Holmes et al.* [2019] by combining the heat
 139 and volume budgets into a budget for *internal heat content* that is independent of the ref-
 140 erence temperature. The internal heat content tendency in *Holmes et al.* [2019] is directly
 141 related to changes in surface fluxes and mixing,

$$\frac{\partial \mathcal{H}_I}{\partial t}(\Theta, t) = \rho_0 C_p \int_{\Theta}^{\infty} \frac{\partial \mathcal{V}}{\partial t}(\Theta, t) d\Theta = \mathcal{F}(\Theta, t) + \mathcal{M}(\Theta, t), \quad (4)$$

142 where $\mathcal{F}(\Theta, t)$ is the surface heat flux into the volume $\mathcal{V}(\Theta, t)$, including the compo-
 143 nent associated with surface volume fluxes (see *Holmes et al.* [2019]). $\mathcal{M}(\Theta, t)$ is the heat
 144 transport across the Θ isotherm due to mixing.

145 OHC change in the fixed-temperature framework results in changes in the volume
 146 distribution of temperature classes (i.e., $d\mathcal{V}/dt$ or equivalently dp/dt). OHC change in
 147 the temperature percentiles results in changes in the temperature of the isotherms that
 148 bound a given volume percentile of the ocean (i.e., $d\Theta^P/dt$). The heat content in the tem-
 149 perature percentiles in equation (3) may thus be related to the OHC in the fixed-temperature
 150 framework in equation (1) using the slope of the cumulative volume distribution of the
 151 ocean in fixed-temperature coordinates (figure 1):

$$\frac{\partial \Theta^P}{\partial t} = \frac{\partial p}{\partial t} \left(\frac{\partial \Theta^P}{\partial p} \right). \quad (5)$$

152 This transformation between fixed-temperature and temperature-percentile co-ordinates
 153 allows us to map between the internal heat content tendency in *Holmes et al.* [2019] and
 154 the equivalent heat content tendency in the tracer-percentile framework. Combining equa-
 155 tions (4) and (5), we obtain the OHC tendency in temperature percentiles:

$$\frac{\partial \mathcal{H}}{\partial t}(p, t) = 0.01 \mathcal{V}_T \rho_0 C_p \int_0^p \frac{\partial \Theta^P}{\partial t}(p, t) dp = \mathcal{F}(p, t) + \mathcal{M}(p, t). \quad (6)$$

156 Equation (6) shows that in temperature percentiles the total heat content tendency is,
 157 like the internal heat content tendency in the fixed-temperature framework [equation (4)],
 158 unaffected by the across-isotherm heat transport associated with across-isotherm volume
 159 transport, as this term is by definition zero. Rather, changes in the global heat content can
 160 be directly attributed to diabatic surface fluxes $\mathcal{F}(p, t)$ and mixing $\mathcal{M}(p, t)$.

161 3 Data

162 3.1 Observations

163 We use a hybrid observational dataset which combines observations of monthly-
 164 averaged, in-situ temperature from two optimally-interpolated gridded datasets:

- 165 1. A temperature field developed by *Cheng and Zhu* [2016] with temperature casts
 166 sourced from the *World Ocean Database*, hereafter referred to as the *Institute of*
 167 *Atmospheric Physics (IAP) dataset*, and
- 168 2. A temperature field from the UK Met Office Hadley Centre Enhanced Ocean Data
 169 Assimilation and Climate prediction (ENACT) version 4 (*EN4*, subversion EN.4.2.1,

170 with *Gouretski and Reseghetti* [2010] corrections, see *Good et al.* [2013] for more
 171 details).

172 The IAP dataset is formulated to reduce sampling errors which arise due to scarce
 173 observations prior to the introduction of the Argo program. This is accomplished by com-
 174 bining observed temperature casts with the error covariance matrix from an ensemble of
 175 CMIP5 models [*Cheng and Zhu*, 2016; *Cheng et al.*, 2017]. The IAP data is only avail-
 176 able for the top 2000m of the ocean, so we fill the deep ocean with temperature data
 177 from EN4. Our hybrid dataset extends from January 1970 to December 2014, and has
 178 a $1^\circ \times 1^\circ$ horizontal grid with 53 vertical levels. We convert in-situ temperature in ob-
 179 servations to conservative temperature as it is proportional to potential enthalpy [*Mc-*
 180 *Dougall*, 2003; *Graham and McDougall*, 2013]. Monthly-averaged variables are binned
 181 in temperature space and time-averaged to yield annual means. Tendency terms are sub-
 182 sequently calculated as the linear trend of a given volume percentile between 1970 and
 183 2014. Marginal seas, namely, the Mediterranean, Red, Baltic, and Black Seas, and the Per-
 184 sian Gulf and Hudson Bay, are excluded from the analysis. Reference density is assumed
 185 to be $\rho_0 = 1035 \text{ kg m}^{-3}$ and the specific heat $C_p = 4000 \text{ J kg}^{-1} \text{ K}^{-1}$.

186 Uncertainty in the observations is estimated by repeating the analysis 1000 times
 187 with temperatures perturbed in a random normal distribution based on standard error
 188 estimates from the IAP dataset. Note that the uncertainty obtained from this method is
 189 smaller than the standard error of the linear fit used to obtain tendencies. Standard error
 190 is not estimated below 1000m in EN4 [*Good et al.*, 2013]. Therefore, in this analysis, un-
 191 certainty is represented by the standard error of the linear regression of the relevant terms.
 192 Note that the auto-correlation coefficient of the OHC tendency is above 0.95 for a time lag
 193 of 1 year at all volume percentiles, so the standard error from the linear regression may be
 194 thought of as a measure of how subsampling different years which have largely indepen-
 195 dent data might affect our analysis.

196 3.2 Models

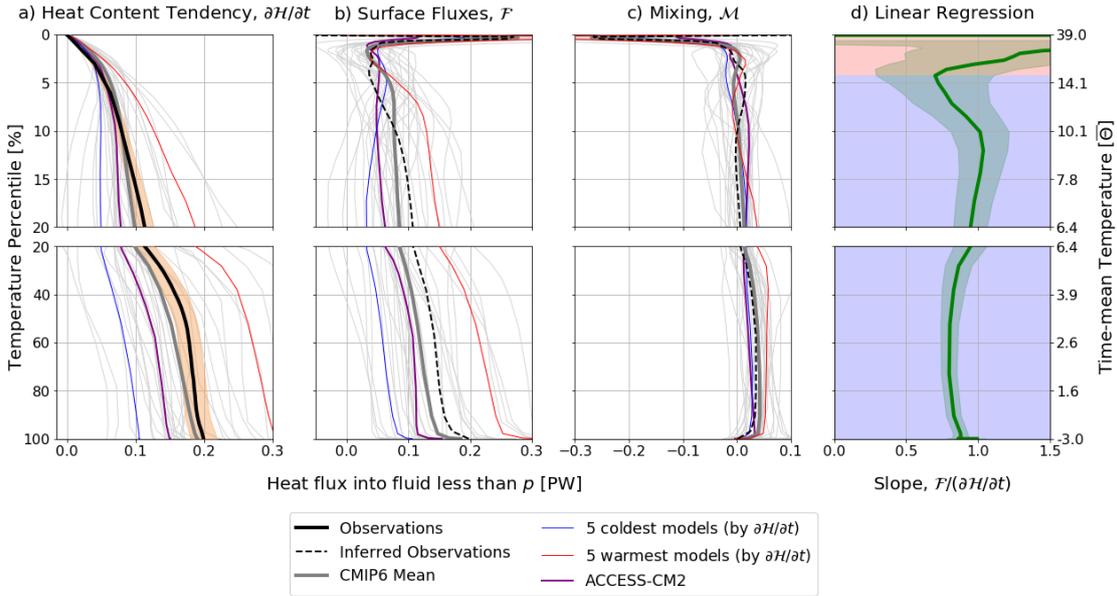
197 Temperature and surface fluxes in thirty climate models that form part of the Cli-
 198 mate Model Intercomparison Project phase 6 [CMIP6; *Eyring et al.*, 2016] are analysed
 199 (see Table S1). We focus on the historical experiment, which branches from the pre-industrial
 200 control (piControl) experiment and covers the time period 1850–2014. The historical runs
 201 include all natural and anthropogenic forcing. The difference between conservative and
 202 potential temperature is negligible, so we use the more widely-reported potential temper-
 203 ature in the CMIP6 calculations. We use monthly-averaged data from the 1970–2014 pe-
 204 riod and mask the Mediterranean, Red, Baltic, and Black Seas, and the Persian Gulf and
 205 Hudson Bay prior to analysis. Model drift is removed in all variables of interest in or-
 206 der to avoid contamination of any forced trends and to ensure closure of the global ocean
 207 heat budget [*Irving et al.*, 2020, see supplementary information]. The models used in this
 208 study archive valid monthly timescale potential temperature data, *thetao*, grid cell vol-
 209 ume, *volcello*, and surface flux data, *hfds* for both the historical and piControl experi-
 210 ments. As with the observations, the monthly-averaged surface fluxes, grid cell volume
 211 and temperature are binned into temperature co-ordinates and time-averaged to yield an-
 212 nual means. The surface flux tendency $\mathcal{F}(\Theta^*, t)$ is calculated as the time-derivative of
 213 the time-integrated annual mean surface flux difference between the historical and pre-
 214 industrial control experiments. More details of this procedure are provided in the supple-
 215 mentary information. The binned diagnostics are then interpolated from fixed-temperature
 216 coordinates to temperature percentiles. The slope of the linear regression of the binned
 217 variables *thetao* and *volcello* in temperature percentiles provides an estimate of $\partial\mathcal{H}/\partial t$
 218 for the models. Details of the individual model members that make up the suite of CMIP6
 219 models analysed are provided in the supplementary information. In the analysis of CMIP6
 220 models, the reference density is $\rho_0 = 1035 \text{ kg m}^{-3}$, and the specific heat is $C_p = 4000 \text{ J kg}^{-1} \text{ K}^{-1}$,

221 unless otherwise specified by modelling groups in the relevant documentation (see refer-
 222 ences in Table S1).

223 To complement our analysis of the standard CMIP6 diagnostics, we assesses a range
 224 of additional outputs from the ACCESS-CM2 historical and piControl simulations which
 225 are not archived as part of CMIP6 [Bi *et al.*, 2020]. The standard surface heat flux diag-
 226 nostic reported for CMIP6 (*hfds*) does not account for the redistribution of shortwave ra-
 227 diation into the ocean interior (which is absorbed at a temperature different to the SST).
 228 To evaluate the impact of this missing process on our results, as well as the role of spe-
 229 cific mixing processes, we use precise ocean heat budget tendency diagnostics from the
 230 ACCESS-CM2 simulations. These tendency diagnostics are binned into fixed-temperature
 231 coordinates using the monthly-averaged temperature distribution.

232 4 Results

233 We quantify the integrated heat content tendency $\partial\mathcal{H}(p,t)/\partial t$ in the observations
 234 and CMIP6 models, the integrated surface flux tendency \mathcal{F} in the CMIP6 models, and the
 235 inferred mixing tendency \mathcal{M} in figure 2. The mixing in figure 2c is the residual between
 236 figures 2a and 2b (with the exception of the ACCESS-CM2 diagnostic term in purple),
 237 and includes any errors associated with binning, the neglect of shortwave penetration and
 238 time-averaging. By construction, the integrated curves in figure 2a, b and c are 0 at $p =$
 239 0% .



240 **Figure 2.** The a) heat content tendency $\partial\mathcal{H}(p,t)/\partial t$, b) surface heat flux trends \mathcal{F} , and c) inferred mixing
 241 changes \mathcal{M} integrated from hot to cold (i.e., values at $p = 100\%$ are the global integral). Tendencies in a)
 242 and b) are calculated as the linear trend in heat content or surface flux at constant temperature-percentile from
 243 1970 to 2014. Orange shading shows the standard error (2σ) of the linear trend in heat content. d) Linear
 244 regression of the integrated \mathcal{F} and $\partial\mathcal{H}/\partial t$ at each temperature-percentile in the CMIP6 models. Green line
 245 indicates the slope of the linear regression and green shading shows the standard error of the linear fit. The
 246 red and blue shading represents the sub-tropical and sub-polar oceans, respectively, based on the time-mean
 247 surface temperature bounds $28^\circ\text{C} > \bar{\Theta} > 15^\circ\text{C}$ and $\bar{\Theta} < 15^\circ\text{C}$, respectively, from Grist *et al.* [2016]. The
 248 tropics (defined as $\bar{\Theta} > 28^\circ\text{C}$) are not visible in this plot. The warmest 20% volume is expanded in these
 249 plots. A secondary y-axis marks the time-mean temperature $\bar{\Theta}$ corresponding to each temperature-percentile.

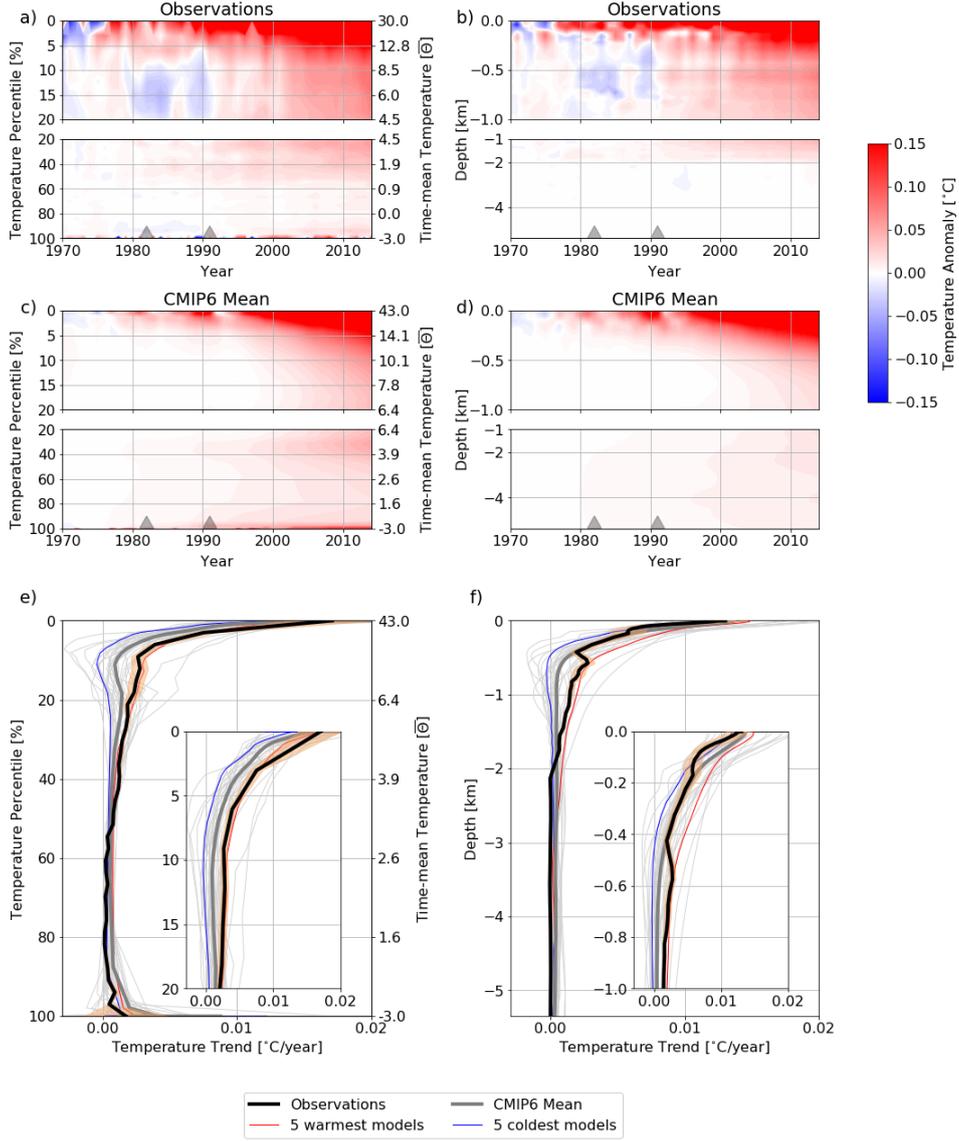
250 The heat content increases across all temperature percentiles in the observations
 251 and the majority of CMIP6 models (figure 2a). There is a reasonable match in OHC ten-
 252 dency between the CMIP6 mean and the observations (compare black and gray lines in
 253 figure 2a), though there is a substantial spread between the models in the CMIP6 ensem-
 254 ble (compare the blue and red lines in figure 2a), and the ensemble-mean OHC is less
 255 than the observed OHC for $p > 10\%$. Surface flux tendencies again show a large spread
 256 between models (figure 2b).

257 The strong increase in surface fluxes in the warmest $\sim 2.5\%$ of the ocean is largely
 258 balanced by mixing, which fluxes this additional heat toward cooler temperature classes.
 259 This region corresponds to warm temperatures ($\Theta > \sim 25^\circ\text{C}$), which collapse to a narrow
 260 temperature-percentile range at the top of the figure. There is also an anomalous diver-
 261 gence of heat due to mixing (driving a cooling tendency from mixing) out of the coldest
 262 40% of waters and an anomalous convergence of heat due to mixing (driving a warming
 263 tendency from mixing) between $10\% < p < 60\%$. The largest cooling tendency from
 264 mixing occurs out of the 5% coldest waters.

265 The surface flux tendencies in figure 2b do not account for shortwave redistribution.
 266 In addition, mixing is not a defined CMIP6 variable, and is not available for many CMIP6
 267 models. Therefore, it remains unclear how much of the the difference between heat con-
 268 tent tendency and surface forcing shown in figure 2c is due to mixing, and how much is
 269 due to shortwave redistribution and other errors. To explore this, we analyse the ocean
 270 model component of a single CMIP6 model member - ACCESS-CM2 - applying the di-
 271 agnostic framework of *Holmes et al.* [2019] to explicitly calculate mixing and the effect of
 272 shortwave redistribution. The breakdown of processes in ACCESS-CM2 (purple lines in
 273 figure 2) largely aligns with the inferred diabatic fluxes in the rest of the CMIP6 models.
 274 There is a cooling tendency from mixing between $60\% < p < 100\%$ and a warming ten-
 275 dency from mixing between $10\% < p < 60\%$. In addition, the strong increase in surface
 276 fluxes at $0\% < p < 2.5\%$ is reduced by 17.7% when shortwave redistribution is taken into
 277 account (not shown). The mixing in these volume classes consequently does not change as
 278 much as inferred from the CMIP6 models.

279 It is not possible to calculate surface flux tendencies and mixing from coarse res-
 280 olution observational datasets. Instead, reanalysis products are often used which rely on
 281 ocean models (e.g., *Liang and Yu* [2016]). The tracer-percentile framework allows us to
 282 infer bulk surface flux and mixing quantities directly from observed OHC tendencies us-
 283 ing the CMIP6 relationships in figure 2. A linear regression analysis is performed be-
 284 tween integrated surface flux tendencies and heat content tendencies at each temperature-
 285 percentile for each of the CMIP6 models, with the slope and error of the regression shown
 286 in figure 2d. The inferred observed surface flux may be calculated using this slope as
 287 $\mathcal{F} = \partial\mathcal{H}/\partial t \times \text{Slope}$. There is significant error in the linear regression in the tropical and
 288 sub-tropical regions (below $p \approx 10\%$). However, at temperature percentiles colder than
 289 $p \approx 10\%$ (i.e. the coldest 90% of the ocean), there is a statistically significant correlation
 290 between heat content and surface flux tendencies. Between $10\% < p < 20\%$, the slope
 291 of the correlation is close to 1, implying that excess heat entering the surface at these vol-
 292 ume classes tends to remain there. By construction, the slope should be 1 at $p = 100\%$,
 293 as there is no mixing in a globally integrated sense. Any deviation from 1 is due to diag-
 294 nosed non-closure of the ocean heat budget in the models [*Irving et al.*, 2020]. Assuming
 295 the relationship between surface fluxes and OHC tendency in the CMIP6 models holds for
 296 the real ocean, we infer the integrated surface flux-driven tendency and mixing-driven ten-
 297 dency as a function of temperature percentiles for the observations (dashed lines in figures
 298 2b and c). Note that the inferred surface flux and mixing profile is only statistically signif-
 299 icant above a temperature-percentile of $\sim 10\%$. A further breakdown of the observed OHC
 300 tendency and its inferred surface fluxes and mixing is discussed in section 5.

309 The OHC tendency in the CMIP6 mean is less than that in the observations (Fig-
 310 ure 2a). The tracer-percentile framework allows us to identify volume classes (organ-



301 **Figure 3.** Global temperature anomaly calculated relative to a 1970-1980 baseline in a) temperature per-
 302 centiles, and b) fixed depth in the observations. c) and d) same as above but for the CMIP6 ensemble-mean.
 303 Grey triangles indicate volcanic eruptions of El Chichón (1982) and Pinatubo (1991). e) and f) Temperature
 304 anomaly trends (°C/year) calculated from the linear trend in temperature at constant temperature-percentile
 305 and depth, respectively. Orange shading shows the standard error (2σ) of the linear trend in temperature
 306 anomaly. The warmest 20% of the ocean (or the top 1000 metres in panels b and d) use an expanded y-axis. A
 307 secondary axis marks the time-mean temperature $\bar{\Theta}$ corresponding to the temperature-percentile in a), c) and
 308 e).

311 ised by temperature) where the cold anomaly that leads to this bias is introduced, and
 312 whether it can be traced back to mixing or surface flux changes. To this end, we plot the
 313 temperature anomaly in temperature percentiles and fixed depths in the observations and
 314 CMIP6 models in figure 3. Recall from section 2 that the temperature anomaly in tem-
 315 perature percentiles is proportional to the percentile derivative of the OHC tendency, i.e.,
 316
$$\frac{\partial \Theta^P}{\partial t} = \frac{100}{\mathcal{V}_T \rho_0 C_p} \frac{\partial}{\partial P} \left(\frac{\partial H}{\partial t} \right).$$

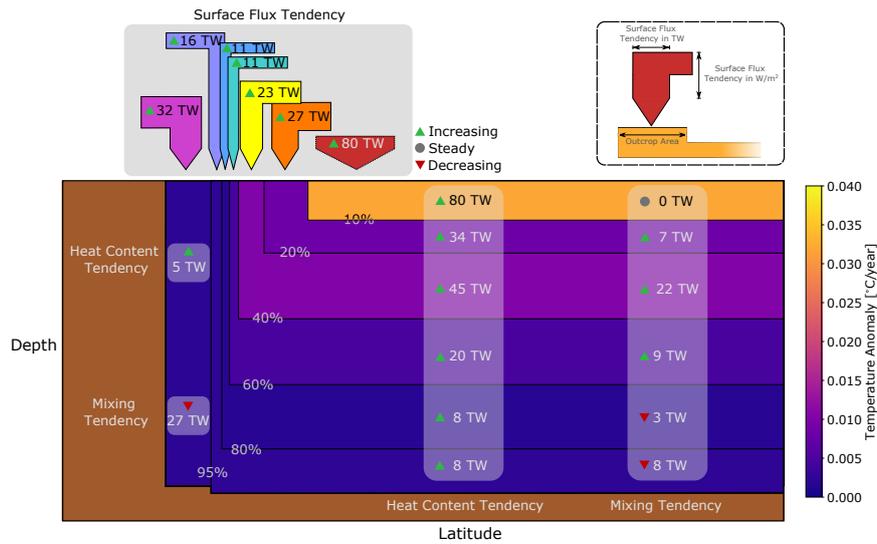
317 In tracer-percentiles temperature anomalies are qualitatively similar to the fixed-
 318 depth framework (figures 3a-d). Note in particular the similarity of the observed negative
 319 temperature anomaly between 1980 and 1990 in tracer-percentiles and fixed-depth coordi-
 320 nates (compare figures 3a and 3b), as well as the persistent cooler (compared to observa-
 321 tions) signal in the CMIP6 mean between $5\% < p < 20\%$ and $-1 \text{ km} < z < -0.25 \text{ km}$
 322 (compare figures 3c and 3d, and thick grey lines in 3e and 3f). A key difference between
 323 the temperature percentiles and fixed depth is the warm anomaly in the coldest 20% tem-
 324 perature percentiles, which cannot be seen in the deep ocean ($z < -4 \text{ km}$). The coldest
 325 percentiles include surface polar waters. Therefore, where warming of surface polar waters
 326 is conflated with changes in the surface sub-tropics and tropics at fixed depth, surface po-
 327 lar warming is emphasised at the coldest temperature percentiles. Analysis of temperature
 328 anomalies of water colder than 0° at fixed-depth (not shown) confirms that the majority
 329 of the warming in the coldest percentiles originates at the surface of the ocean. The tem-
 330 perature trend over the full time period, shown in temperature-percentile and fixed-depth
 331 coordinates in figures 3e and 3f respectively, reveals that the observed annual temperature
 332 anomaly has a number of ‘kinks’ in the sub-surface ($z \approx [0.2, 0.5] \text{ km}$) in fixed-depth co-
 333 ordinates which are not visible in the tracer-percentile framework. We posit these kinks
 334 may be associated with XBT corrections in the observations [Abraham *et al.*, 2013] or in-
 335 ternal variability.

336 The CMIP6 models exhibit a weaker warming trend in the $\sim 40\%$ warmest vol-
 337 ume classes in the ocean, particularly between $5\% < p < 20\%$ (compare black and gray
 338 lines in figure 3e). Revisiting the CMIP6 linear regression in figure 2d, we conclude that
 339 the weaker modelled warming trend in volume classes above $\sim 10\%$ is likely attributed
 340 to biases in surface flux changes (as the regression coefficient is close to 1 at these per-
 341 centiles). These surface flux trends may be inaccurate in magnitude or spatial pattern,
 342 leading to heat entering the ocean at incorrect volume classes. Past research with CMIP5
 343 models has also identified a consistent sea surface temperature (SST) bias in climate mod-
 344 els - particularly a warm bias in the Southern Ocean SSTs [Sallée *et al.*, 2013; Meijers,
 345 2014]. There is evidence this SST bias still exists in the current generation of CMIP6
 346 models [Beadling *et al.*, 2020]. In fact, SSTs and surface fluxes are tightly coupled, so
 347 the true reason for the CMIP6 bias is likely to be a combination of inaccuracies in both
 348 [Hyder *et al.*, 2018]. A possible explanation for the weaker CMIP6 warming trend may
 349 also be that the CMIP6 models have biased mixing, which would imply a surface flux bias
 350 elsewhere.

351 5 Discussion and Conclusions

352 We introduce a novel diagnostic framework (termed the *tracer-percentile frame-*
 353 *work*) which we apply to ocean temperature and use to directly relate changes in OHC
 354 to changes in diabatic heat fluxes, namely, surface forcing and mixing. By quantifying
 355 the heat content in and transport across temperature percentiles (ordered from warmest to
 356 coldest), we exclude adiabatic processes and avoid regime changes associated with shifting
 357 isotherms in a warming ocean. We are thus able to trace changes in OHC to changes in
 358 surface fluxes and mixing using a combination of traditional observational hydrographic
 359 datasets and CMIP6 models. Comparing heat content tendency and surface fluxes in a
 360 suite of thirty CMIP6 models from 1970 to 2014, we establish that changes in the cold-
 361 est 90% of the ocean may be traced to changes in the net surface heat flux into volume
 362 classes in the sub-polar ocean. Assuming that the linear regression between OHC ten-
 363 dency and surface flux changes in the CMIP6 climate models holds for the real ocean, we
 364 infer changes in sub-polar surface heat fluxes and mixing based on observed OHC trends.

375 Figure 4 summarises the observed OHC tendency and the inferred surface flux and
 376 mixing tendencies in temperature-percentile layers in the ocean. We find that $\sim 60\%$ of
 377 the anomalous surface flux into the ocean (sum of all arrows in figure 4 save the red ar-
 378 row) enters 23% of the surface ocean (sum of all outcrop lengths in figure 4 except the



365 **Figure 4.** Observed temperature anomalies, heat content tendencies, and inferred surface flux and mix-
 366 ing tendencies using the tracer-percentile framework. The OHC tendency, surface fluxes and mixing are
 367 calculated as the derivative of the solid and dashed black curves in figures 2 a, b and c. Note that while the
 368 Southern Hemisphere is schematised here, the values shown are for the global ocean. By construction, mixing
 369 tendencies sum to zero globally, and a decrease in down-gradient mixing implies a reduction in OHC ten-
 370 dency in that layer, and vice versa. The area of each layer in the schematic is proportional to its volume, and
 371 the outcrop length at the top surface is proportional to the outcrop area of the layer. The width of the surface
 372 flux arrows is proportional to their value (in TW), and the height is proportional to the surface flux per unit
 373 unit area (in W/m^2). Surface flux and mixing trends in the 10% warmest ocean volume are uncertain due to the
 374 poor linear fit in CMIP6 models.

379 outcrop of the warmest 10% volume), which represents the outcrop area of the temperature-
380 percentile bounding 90% of the ocean by volume. Therefore, minor changes to SSTs or
381 surface fluxes overlying this relatively small surface area of the ocean may have profound
382 impacts on the mean stratification and ocean circulation. Down-gradient mixing is de-
383 creasing across the 40%-95% percentiles, and increasing across the 10% and 20% per-
384 centiles according to our analysis. The strongest decrease in mixing occurs across the 95%
385 percentile.

386 For the first time, the tracer-percentile framework allows us to link observed OHC
387 changes to surface flux tendencies and the surface area over which they enter. This result
388 enables the identification of the cause of biases in OHC tendencies in climate models. We
389 identify a cooling bias in CMIP6 models which is apparent at both fixed depth and tem-
390 perature percentiles. We trace this cooling bias to anomalous changes in surface forcing
391 into the same temperature percentiles. The surface flux field may have either incorrect
392 magnitude or incorrect spatial distribution, the SSTs which define the volume classes into
393 which the surface fluxes enter may be biased in many ocean models, or the mixing in all
394 CMIP6 models may be biased (with a surface flux bias elsewhere).

395 The insights from the tracer-percentile framework motivate a mechanistic approach
396 to future explorations of OHC tendency and its diabatic contributors. The global analy-
397 sis presented here indicates that surface fluxes over a small region of the ocean contribute
398 to OHC changes in the vast majority of the ocean volume. Therefore, understanding the
399 mechanisms which lead to increases in surface fluxes in these regions (and indeed to sur-
400 face flux biases in models) is crucial. Mixing also displays variability between volume
401 classes, and understanding the processes which drive these changes in mixing is vital to
402 gaining a complete picture of historical and future changes in OHC. The tracer-percentile
403 framework is a versatile tool, and future efforts will work towards expanding the use-cases
404 of the tracer-percentile framework and further understanding oceanic and atmospheric dy-
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415 data/ocean-temperature-analysis-and-heat-content-estimate-institute-atmospheric-physics](https://climatedataguide.ucar.edu/climate-data/ocean-temperature-analysis-and-heat-content-estimate-institute-atmospheric-physics)).
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